

Lecture 17 - Supplement: Derivation of Navier Stokes

* Navier-stokes is a force balance written in a slightly different way: as a momentum balance.

* See also
Deen 6.4, 6.6

$$\sum_i \underline{F}_i = m \underline{a} \quad \underline{a} = \frac{d\underline{v}}{dt}$$

↑ if mass is constant
then we can pull it into $\frac{d}{dt}$

$$\sum_i \underline{F}_i = m \frac{d\underline{v}}{dt} = \frac{d}{dt} (m \underline{v})$$

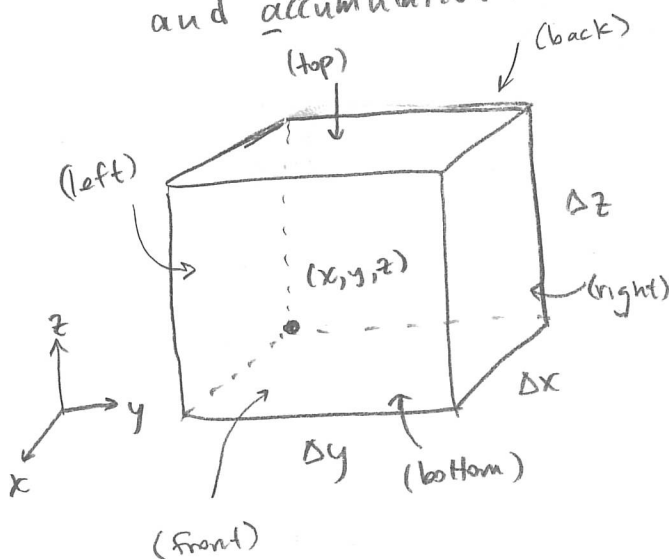
↑ momentum, $\underline{p}$

$$\sum_i \underline{F}_i = \frac{d\underline{p}}{dt}$$

↑
generation terms accumulation term

what generates momentum?
Forces!

* A balance equation will have input, output, generation and accumulation terms: IOGA .



$$\text{accum} = \text{in} - \text{out} + \text{generation}$$

$$\frac{d}{dt}(\underline{p}) = \underbrace{\dot{p}_{\text{in}} - \dot{p}_{\text{out}}}_{\text{momentum flow}} + \underbrace{\sum_i \underline{F}_i}_{\text{sum of forces on C.V.}}$$

↑ control volume

⊗ Notice this is a vector equation!

A. Accumulation term.

$$\frac{\partial}{\partial t}(\underline{p}) = \frac{\partial}{\partial t}(m \underline{v}) = \frac{\partial}{\partial t} \left(\underbrace{\rho \Delta x \Delta y \Delta z}_{\text{box volume}} \underline{v} \right)$$

↑
partial derivative
b/c we will have other space derivatives

B. In-Out terms.

what is the momentum flow: $\dot{\underline{p}}_{\text{in}} - \dot{\underline{p}}_{\text{out}}?$

• recall mass flow: $\dot{m} = \rho Q = \rho u A = \rho \underline{v} \cdot (-\underline{n} A)$

mass flux = $\frac{\text{mass}}{\text{area time}}$

negative b/c we want "in" to be positive

• by analogy we can write for momentum:

$$\dot{\underline{p}} = \dot{m} \underline{v} = \rho Q \underline{v} = \rho u \underline{v} A$$

$$= \underbrace{\rho \underline{v} \underline{v}} \cdot (-\underline{n} A)$$

momentum flux = $\frac{\text{momentum}}{\text{area time}}$

dyadic

$$\rho \underline{v} \underline{v} = \rho \begin{bmatrix} v_x v_x & v_x v_y & v_x v_z \\ v_y v_x & v_y v_y & v_y v_z \\ v_z v_x & v_z v_y & v_z v_z \end{bmatrix}$$

* Now we need to evaluate the momentum flow on all 6 sides of our control volume:

(back) $-(\rho \underline{v} \underline{v})|_x \cdot (-\underline{e}_x) \Delta y \Delta z = (\rho \underline{v} v_x)|_x \Delta y \Delta z$

(front) $-(\rho \underline{v} \underline{v})|_{x+\Delta x} \cdot (\underline{e}_x) \Delta y \Delta z = -(\rho \underline{v} v_x)|_{x+\Delta x} \Delta y \Delta z$

$$(left) \quad -(\rho \underline{u} \underline{u})|_y \cdot (-\underline{e}_y) \Delta x \Delta z = (\rho \underline{u} u_y)|_y \Delta x \Delta z$$

$$(right) \quad -(\rho \underline{u} \underline{u})|_{y+\Delta y} \cdot (\underline{e}_y) \Delta x \Delta z = -(\rho \underline{u} u_y)|_{y+\Delta y} \Delta x \Delta z$$

$$(bottom) \quad -(\rho \underline{u} \underline{u})|_z \cdot (-\underline{e}_z) \Delta x \Delta y = (\rho \underline{u} u_z)|_z \Delta x \Delta y$$

$$(top) \quad -(\rho \underline{u} \underline{u})|_{z+\Delta z} \cdot (\underline{e}_z) \Delta x \Delta y = -(\rho \underline{u} u_z)|_{z+\Delta z} \Delta x \Delta y$$

* all together:

$$\begin{aligned} \dot{P}_{in} - \dot{P}_{out} = & \left[(\rho \underline{u} u_x)|_x - (\rho \underline{u} u_x)|_{x+\Delta x} \right] \Delta y \Delta z \\ & + \left[(\rho \underline{u} u_y)|_y - (\rho \underline{u} u_y)|_{y+\Delta y} \right] \Delta x \Delta z \\ & + \left[(\rho \underline{u} u_z)|_z - (\rho \underline{u} u_z)|_{z+\Delta z} \right] \Delta x \Delta y \end{aligned}$$

C. Generation

* Now we need to look at the sum of the forces on our control volume:

$$\sum_i \underline{F} = \underline{F}_s + \underline{F}_b$$

$\uparrow$ $\uparrow$
 surface body
 forces forces

$$\underline{F}_b = m \underline{g} \quad \leftarrow \text{gravitational forces}$$

$$\underline{F}_s = A \underline{n} \cdot \underline{\sigma} = A \left(\underline{n} P + \underline{n} \cdot \underline{\tau} \right)$$

$\uparrow$ $\uparrow$
 pressure viscous stresses

Body Forces: $\underline{F}_b = m \underline{g} = \rho \Delta x \Delta y \Delta z \underline{g}$

Surface forces: we have 6 surfaces!

(back) $-\underline{e}_x \cdot \underline{\sigma}|_x \Delta y \Delta z$

(front) $\underline{e}_x \cdot \underline{\sigma}|_{x+\Delta x} \Delta y \Delta z$

(left) $-\underline{e}_y \cdot \underline{\sigma}|_y \Delta x \Delta z$

(right) $\underline{e}_y \cdot \underline{\sigma}|_{y+\Delta y} \Delta x \Delta z$

(bottom) $-\underline{e}_z \cdot \underline{\sigma}|_z \Delta x \Delta y$

(top) $\underline{e}_z \cdot \underline{\sigma}|_{z+\Delta z} \Delta x \Delta y$

* combining the forces gives:

$$\begin{aligned} \sum_i \underline{F}_i &= \rho \Delta x \Delta y \Delta z \underline{g} + \underline{e}_x \cdot (\underline{\sigma}|_{x+\Delta x} - \underline{\sigma}|_x) \Delta y \Delta z \\ &\quad + \underline{e}_y (\underline{\sigma}|_{y+\Delta y} - \underline{\sigma}|_y) \Delta x \Delta z \\ &\quad + \underline{e}_z (\underline{\sigma}|_{z+\Delta z} - \underline{\sigma}|_z) \Delta x \Delta y \end{aligned}$$

D. Cauchy Momentum Equation

$$\frac{\partial}{\partial t}(\underline{p}) = \underline{\dot{p}}_{in} - \underline{\dot{p}}_{out} + \sum_i \underline{F}_i$$

combining all of our previous pieces together

we get:

$$\begin{aligned}
\frac{\partial}{\partial t}(\rho \Delta x \Delta y \Delta z \underline{v}) = & - \left[(\rho \underline{v} v_x)|_{x+\Delta x} - (\rho \underline{v} v_x)|_x \right] \Delta y \Delta z \\
& - \left[(\rho \underline{v} v_y)|_{y+\Delta y} - (\rho \underline{v} v_y)|_y \right] \Delta x \Delta z \\
& - \left[(\rho \underline{v} v_z)|_{z+\Delta z} - (\rho \underline{v} v_z)|_z \right] \Delta x \Delta y \\
& + \underline{e}_x \cdot (\underline{\sigma}|_{x+\Delta x} - \underline{\sigma}|_x) \Delta y \Delta z \\
& + \underline{e}_y \cdot (\underline{\sigma}|_{y+\Delta y} - \underline{\sigma}|_y) \Delta x \Delta z \\
& + \underline{e}_z \cdot (\underline{\sigma}|_{z+\Delta z} - \underline{\sigma}|_z) \Delta x \Delta y \\
& + \rho \Delta x \Delta y \Delta z \underline{g}
\end{aligned}$$

- divide by $\Delta x \Delta y \Delta z$
- take the limit as $\Delta x \rightarrow 0, \Delta y \rightarrow 0, \Delta z \rightarrow 0$

$$\frac{(\rho \underline{v} v_x)|_{x+\Delta x} - (\rho \underline{v} v_x)|_x}{\Delta x} \rightarrow \frac{\partial}{\partial x}(\rho \underline{v} v_x)$$

and so on...

$$\frac{\partial}{\partial t}(\rho \underline{v}) = - \frac{\partial}{\partial x}(\rho \underline{v} v_x) - \frac{\partial}{\partial y}(\rho \underline{v} v_y) - \frac{\partial}{\partial z}(\rho \underline{v} v_z)$$

$$+ \underline{e}_x \cdot \frac{\partial \underline{\sigma}}{\partial x} + \underline{e}_y \cdot \frac{\partial \underline{\sigma}}{\partial y} + \underline{e}_z \cdot \frac{\partial \underline{\sigma}}{\partial z} + \rho \underline{g}$$

Can write these

guys more compactly: $\nabla \cdot \underline{f} = \underline{e}_x \cdot \frac{\partial \underline{f}}{\partial x} + \underline{e}_y \cdot \frac{\partial \underline{f}}{\partial y} + \underline{e}_z \cdot \frac{\partial \underline{f}}{\partial z}$

$$\frac{\partial}{\partial t}(\rho \underline{v}) = - \nabla \cdot (\rho \underline{v} \underline{v}) + \nabla \cdot \underline{\sigma} + \rho \underline{g} \quad (\text{phew!})$$

$$\underbrace{\frac{\partial}{\partial t}(\rho \underline{v}) + \nabla \cdot (\rho \underline{v} \underline{v})}_{\text{this term will simplify}} = \nabla \cdot \underline{\sigma} + \rho \underline{g}$$

↖ this term will simplify

$$\frac{\partial}{\partial t}(\rho \underline{v}) = \rho \frac{\partial \underline{v}}{\partial t} + \underline{v} \frac{\partial \rho}{\partial t}$$

$$\nabla \cdot (\rho \underline{v} \underline{v}) = \rho \underline{v} \cdot \nabla \underline{v} + \underline{v} \nabla \cdot (\rho \underline{v})$$

← This is the product rule in vector calculus
 (*) See appendix

$$\begin{aligned} \frac{\partial}{\partial t}(\rho \underline{v}) + \nabla \cdot (\rho \underline{v} \underline{v}) &= \rho \frac{\partial \underline{v}}{\partial t} + \underline{v} \frac{\partial \rho}{\partial t} + \rho \underline{v} \cdot \nabla \underline{v} + \underline{v} \nabla \cdot (\rho \underline{v}) \\ &= \underline{v} \left[\underbrace{\frac{\partial \rho}{\partial t} + \nabla \cdot (\rho \underline{v})}_{\text{continuity equation}} \right] + \rho \left[\underbrace{\frac{\partial \underline{v}}{\partial t} + \underline{v} \cdot \nabla \underline{v}}_{\text{material derivative}} \right] \end{aligned}$$

$$\frac{\partial \rho}{\partial t} + \nabla \cdot (\rho \underline{v}) = 0$$

$$\rho \frac{D \underline{v}}{Dt}$$

$$\rho \frac{D \underline{v}}{Dt} = \nabla \cdot \underline{\sigma} + \rho \underline{g}$$

Cauchy's momentum

Equation, "Newton's 2nd law for fluids"

E. Navier - Stokes Equation

* 2 assumptions :

- (1) Newtonian Fluid (constant μ)
- (2) Incompressible fluid (constant ρ)

* Both of these will let us simplify $\underline{\sigma}$:

$$\nabla \cdot \underline{\sigma} = \nabla \cdot (-P \underline{\delta} + \underline{\tau})$$

← Newtonian Fluid

$$\underline{\tau} = 2\mu \underline{\Gamma} = \mu [(\nabla \underline{v}) + (\nabla \underline{v})^T]$$

← also assumes incompressible.

$$\underline{\nabla} \cdot \underline{\underline{\sigma}} = \underline{\nabla} \cdot \{ -P \underline{\underline{\delta}} + \mu [(\underline{\nabla} \underline{v}) + (\underline{\nabla} \underline{v})^T] \}$$

$$= -\underline{\nabla} P + \mu \underline{\nabla} \cdot [(\underline{\nabla} \underline{v}) + (\underline{\nabla} \underline{v})^T]$$

⊛ see appendix
for proof of
vector
property.

$$= -\underline{\nabla} P + \mu \underline{\nabla} \cdot (\underline{\nabla} \underline{v}) + \mu \underline{\nabla} (\underline{\nabla} \cdot \underline{v})$$

$$\underbrace{\underline{\nabla} \cdot \underline{\nabla}} = \nabla^2$$

↑
Incompressible fluid

$$\underline{\nabla} \cdot \underline{v} = 0$$

$$= -\underline{\nabla} P + \mu \nabla^2 \underline{v}$$

Finally,

$$\rho \frac{D \underline{v}}{Dt} = -\underline{\nabla} P + \mu \nabla^2 \underline{v} + \rho \underline{g}$$

← The Navier-Stokes Equation!

- Newton's 2nd law
- + Newtonian
- + Incompressible.

✓ in Cartesian coordinates (p. 147 in §6.6)

$$\rho \left[\frac{\partial v_x}{\partial t} + v_x \frac{\partial v_x}{\partial x} + v_y \frac{\partial v_x}{\partial y} + v_z \frac{\partial v_x}{\partial z} \right] = -\frac{\partial P}{\partial x} + \mu \left[\frac{\partial^2 v_x}{\partial x^2} + \frac{\partial^2 v_x}{\partial y^2} + \frac{\partial^2 v_x}{\partial z^2} \right] + \rho g_x$$

$$\rho \left[\frac{\partial v_y}{\partial t} + v_x \frac{\partial v_y}{\partial x} + v_y \frac{\partial v_y}{\partial y} + v_z \frac{\partial v_y}{\partial z} \right] = -\frac{\partial P}{\partial y} + \mu \left[\frac{\partial^2 v_y}{\partial x^2} + \frac{\partial^2 v_y}{\partial y^2} + \frac{\partial^2 v_y}{\partial z^2} \right] + \rho g_y$$

$$\rho \left[\frac{\partial v_z}{\partial t} + v_x \frac{\partial v_z}{\partial x} + v_y \frac{\partial v_z}{\partial y} + v_z \frac{\partial v_z}{\partial z} \right] = -\frac{\partial P}{\partial z} + \mu \left[\frac{\partial^2 v_z}{\partial x^2} + \frac{\partial^2 v_z}{\partial y^2} + \frac{\partial^2 v_z}{\partial z^2} \right] + \rho g_z$$

F. Appendix : proof of vector identities

$$\begin{aligned}
 \nabla \cdot (\rho \underline{v}) &= \partial_i (\rho v_i v_j) \quad \swarrow \text{Einstein Index Notation} \\
 &= v_j \partial_i (\rho v_i) + \rho v_i \partial_i (v_j) \quad \searrow \text{product rule} \\
 &= \underline{v} \cdot \nabla (\rho \underline{v}) + \rho \underline{v} \cdot \nabla \underline{v}
 \end{aligned}$$

$$\begin{aligned}
 \nabla \cdot [(\nabla \underline{v}) + (\nabla \underline{v})^T] &= \partial_i [\partial_i v_j + \partial_j v_i] \\
 &= \partial_i \partial_i v_j + \partial_i \partial_j v_i \\
 &= \partial_i \partial_i v_j + \partial_j \partial_i v_i \\
 &= \nabla \cdot \nabla (\underline{v}) + \nabla (\nabla \cdot \underline{v}) \\
 &= \nabla^2 \underline{v} + \nabla (\nabla \cdot \underline{v})
 \end{aligned}$$